

Fluid Dynamics Equation Sheet

1 Conservation Equations

1.1 Integral Form

1.1.1 Mass

$$\frac{\partial}{\partial t} \iiint \rho dV + \iint \rho \mathbf{u} \cdot d\mathbf{S} = 0$$

1.1.2 Momentum

$$\begin{aligned} \frac{\partial}{\partial t} \iiint \rho \mathbf{u} dV + \iint \rho (\mathbf{u} \cdot d\mathbf{S}) \mathbf{u} \\ = \iiint \rho \mathbf{F} dV - \iint p d\mathbf{S} + \mathbf{F}_{viscous} \end{aligned}$$

1.1.3 Energy

$$\begin{aligned} \frac{\partial}{\partial t} \iiint \rho \left(e + \frac{1}{2} \mathbf{u} \cdot \mathbf{u} \right) dV + \iint \rho \left(e + \frac{1}{2} \mathbf{u} \cdot \mathbf{u} \right) \mathbf{u} \cdot d\mathbf{S} \\ = \iiint \dot{q} \rho dV + \dot{Q}_{viscous} - \iint p \mathbf{u} \cdot d\mathbf{S} \\ + \iiint \rho (\mathbf{f} \cdot \mathbf{u}) dV + \dot{W}_{viscous} \end{aligned}$$

1.2 Differential Form

1.2.1 Mass

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{u}) = 0$$

1.2.2 Momentum

$$\begin{aligned} \rho \left(\frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} \right) &= \rho \mathbf{F} - \nabla p + \nabla \cdot \boldsymbol{\tau} \\ &= \rho \mathbf{F} - \nabla p + \mu \nabla^2 \mathbf{u} + \frac{1}{3} \mu \nabla (\nabla \cdot \mathbf{u}) \\ &\quad + \kappa \nabla (\nabla \cdot \mathbf{u}) + 2 (\nabla \mu) \cdot \nabla \mathbf{u} + (\nabla \mu) \times (\nabla \times \mathbf{u}) \\ &\quad - \frac{2}{3} (\nabla \mu) (\nabla \cdot \mathbf{u}) + (\nabla \kappa) (\nabla \cdot \mathbf{u}) \end{aligned}$$

$$\kappa = \lambda + \frac{2}{3} \mu \quad (\text{bulk viscosity})$$

$$\boldsymbol{\tau} = \mu [\nabla \mathbf{u} + (\nabla \mathbf{u})^T] + \left(\kappa - \frac{2}{3} \mu \right) (\nabla \cdot \mathbf{u}) \mathbf{I}$$

$$\mathbf{I} = \sum_{i=1}^3 \sum_{j=1}^3 \delta_{ij}$$

1.2.3 Energy

$$\rho \frac{De}{Dt} = -p (\nabla \cdot \mathbf{u}) + \nabla \cdot (k \nabla T) + \boldsymbol{\tau} : \nabla \mathbf{u}$$

$$\boldsymbol{\tau} : \nabla \mathbf{u} = \sum_{i=1}^3 \sum_{j=1}^3 \tau_{ij} (\nabla \mathbf{u})_{ij}$$

1.2.4 Vorticity Transport

$$\begin{aligned} \frac{\partial \boldsymbol{\Omega}}{\partial t} + \nabla \times (\boldsymbol{\Omega} \times \mathbf{u}) &= \nabla \times \mathbf{F} + \frac{\nabla \rho \times \nabla p}{\rho^2} \\ &\quad + \nu \nabla^2 \boldsymbol{\Omega} + (\nabla \nu) \times \left[\frac{3}{2} \nabla (\nabla \cdot \mathbf{u}) - \nabla^2 \boldsymbol{\Omega} \right] \end{aligned}$$

$$\nabla \times (\boldsymbol{\Omega} \times \mathbf{u}) = \boldsymbol{\Omega} (\nabla \cdot \mathbf{u}) + (\mathbf{u} \cdot \nabla) \boldsymbol{\Omega} - (\boldsymbol{\Omega} \cdot \nabla) \mathbf{u}$$

$$\boldsymbol{\Omega} = \nabla \times \mathbf{u} \quad \text{and} \quad \nu = \mu / \rho$$

2 Vector Calculus Identities

2.1 Properties of vector product

$$\begin{aligned}\mathbf{A} \times \mathbf{B} &= -\mathbf{B} \times \mathbf{A} \\ \mathbf{A} \times (\mathbf{B} \times \mathbf{C}) &= -\mathbf{A} \times (\mathbf{C} \times \mathbf{B}) = -(\mathbf{B} \times \mathbf{C}) \times \mathbf{A}\end{aligned}$$

2.2 Differentiation of vectors

$$\begin{aligned}\frac{d\mathbf{A}}{dt} &= \frac{dA_1}{dt}\mathbf{i} + \frac{dA_2}{dt}\mathbf{j} + \frac{dA_3}{dt}\mathbf{k} \\ \frac{d}{dt}(\mathbf{A} + \mathbf{B}) &= \frac{d\mathbf{A}}{dt} + \frac{d\mathbf{B}}{dt} \\ \frac{d}{dt}(\phi\mathbf{A}) &= \frac{d\phi}{dt}\mathbf{A} + \phi\frac{d\mathbf{A}}{dt} \\ \frac{d}{dt}(\mathbf{A} \cdot \mathbf{B}) &= \frac{d\mathbf{A}}{dt} \cdot \mathbf{B} + \mathbf{A} \cdot \frac{d\mathbf{B}}{dt} \\ \frac{d}{dt}(\mathbf{A} \times \mathbf{B}) &= \frac{d\mathbf{A}}{dt} \times \mathbf{B} + \mathbf{A} \times \frac{d\mathbf{B}}{dt} \\ \frac{d}{dt}(\mathbf{A} \times \mathbf{B} \cdot \mathbf{C}) &= \frac{d\mathbf{A}}{dt} \times \mathbf{B} \cdot \mathbf{C} + \mathbf{A} \times \frac{d\mathbf{B}}{dt} \cdot \mathbf{C} + \mathbf{A} \times \mathbf{B} \cdot \frac{d\mathbf{C}}{dt}\end{aligned}$$

2.3 First Derivative Identities

2.3.1 Distributive Properties

$$\begin{aligned}\nabla(\psi + \phi) &= \nabla\psi + \nabla\phi \\ \nabla(\mathbf{A} + \mathbf{B}) &= \nabla\mathbf{A} + \nabla\mathbf{B} \\ \nabla \cdot (\mathbf{A} + \mathbf{B}) &= \nabla \cdot \mathbf{A} + \nabla \cdot \mathbf{B} \\ \nabla \times (\mathbf{A} + \mathbf{B}) &= \nabla \times \mathbf{A} + \nabla \times \mathbf{B}\end{aligned}$$

2.3.2 Product Rule for Multiplication by a Scalar

$$\begin{aligned}\nabla(\psi\phi) &= \phi\nabla\psi + \psi\nabla\phi \\ \nabla(\psi\mathbf{A}) &= (\nabla\psi)\mathbf{A}^\top + \psi\nabla\mathbf{A} = \nabla\psi \otimes \mathbf{A} + \psi\nabla\mathbf{A} \\ \nabla \cdot (\psi\mathbf{A}) &= \psi\nabla \cdot \mathbf{A} + (\nabla\psi) \cdot \mathbf{A} \\ \nabla \times (\psi\mathbf{A}) &= \psi\nabla \times \mathbf{A} + (\nabla\psi) \times \mathbf{A} \\ \nabla^2(fg) &= f\nabla^2g + 2\nabla f \cdot \nabla g + g\nabla^2f\end{aligned}$$

2.3.3 Quotient rule for division by a scalar

$$\begin{aligned}\nabla\left(\frac{\psi}{\phi}\right) &= \frac{\phi\nabla\psi - \psi\nabla\phi}{\phi^2} \\ \nabla\left(\frac{\mathbf{A}}{\phi}\right) &= \frac{\phi\nabla\mathbf{A} - \nabla\phi \otimes \mathbf{A}}{\phi^2} \\ \nabla \cdot \left(\frac{\mathbf{A}}{\phi}\right) &= \frac{\phi\nabla \cdot \mathbf{A} - \nabla\phi \cdot \mathbf{A}}{\phi^2} \\ \nabla \times \left(\frac{\mathbf{A}}{\phi}\right) &= \frac{\phi\nabla \times \mathbf{A} - \nabla\phi \times \mathbf{A}}{\phi^2}\end{aligned}$$

2.3.4 Dot Product Rule

$$\begin{aligned}\nabla(\mathbf{A} \cdot \mathbf{B}) &= (\mathbf{A} \cdot \nabla)\mathbf{B} + (\mathbf{B} \cdot \nabla)\mathbf{A} + \mathbf{A} \times (\nabla \times \mathbf{B}) + \mathbf{B} \times (\nabla \times \mathbf{A}) \\ \frac{1}{2}\nabla(\mathbf{A} \cdot \mathbf{A}) &= (\mathbf{A} \cdot \nabla)\mathbf{A} + \mathbf{A} \times (\nabla \times \mathbf{A})\end{aligned}$$

2.3.5 Cross Product Rule

$$\begin{aligned}\nabla \cdot (\mathbf{A} \times \mathbf{B}) &= (\nabla \times \mathbf{A}) \cdot \mathbf{B} - \mathbf{A} \cdot (\nabla \times \mathbf{B}) \\ \nabla \times (\mathbf{A} \times \mathbf{B}) &= \mathbf{A}(\nabla \cdot \mathbf{B}) - \mathbf{B}(\nabla \cdot \mathbf{A}) + (\mathbf{B} \cdot \nabla)\mathbf{A} - (\mathbf{A} \cdot \nabla)\mathbf{B} \\ &= \nabla \cdot (\mathbf{B}\mathbf{A}^\top - \mathbf{A}\mathbf{B}^\top) \\ \mathbf{A} \times (\nabla \times \mathbf{B}) &= (\nabla\mathbf{B}) \cdot \mathbf{A} - (\mathbf{A} \cdot \nabla)\mathbf{B} \\ (\mathbf{A} \times \nabla) \times \mathbf{B} &= (\nabla\mathbf{B}) \cdot \mathbf{A} - \mathbf{A}(\nabla \cdot \mathbf{B}) \\ &= \mathbf{A} \times (\nabla \times \mathbf{B}) + (\mathbf{A} \cdot \nabla)\mathbf{B} - \mathbf{A}(\nabla \cdot \mathbf{B})\end{aligned}$$

2.4 Second Derivative Identities

2.4.1 Divergence of curl is zero

$$\nabla \cdot (\nabla \times \mathbf{A}) = 0$$

2.4.2 Divergence of gradient is Laplacian

$$\Delta\psi = \nabla^2\psi = \nabla \cdot (\nabla\psi)$$

2.4.3 Divergence of divergence is not defined

$$\nabla \cdot (\nabla \cdot \mathbf{A}) \text{ is undefined}$$

2.4.4 Curl of gradient is zero

$$\nabla \times (\nabla\phi) = 0$$

2.4.5 Curl of curl

$$\nabla \times (\nabla \times \mathbf{A}) = \nabla(\nabla \cdot \mathbf{A}) - \nabla^2\mathbf{A}$$

Here ∇^2 is the vector Laplacian operating on the vector field \mathbf{A} .

2.4.6 Curl of divergence is not defined

$$\nabla \times (\nabla \cdot \mathbf{A}) \text{ is undefined}$$

2.5 Third Derivative Identities

$$\begin{aligned}\nabla^2(\nabla\psi) &= \nabla[\nabla \cdot (\nabla\psi)] = \nabla(\nabla^2\psi) \\ \nabla^2(\nabla \cdot \mathbf{A}) &= \nabla \cdot [\nabla(\nabla \cdot \mathbf{A})] = \nabla \cdot (\nabla^2\mathbf{A}) \\ \nabla^2(\nabla \times \mathbf{A}) &= -\nabla \times [\nabla \times (\nabla \times \mathbf{A})] = \nabla \times (\nabla^2\mathbf{A})\end{aligned}$$

3 Vector Differential Invariants

$$\nabla\psi \equiv \text{grad } \psi = \frac{\mathbf{i}_i}{h_i} \frac{\partial\psi}{\partial q_i}$$

$$\nabla \cdot \mathbf{u} \equiv \text{div } \mathbf{u} = \frac{1}{h_1 h_2 h_3} \left[\frac{\partial(u_1 h_2 h_3)}{\partial q_1} + \frac{\partial(u_2 h_3 h_1)}{\partial q_2} + \frac{\partial(u_3 h_1 h_2)}{\partial q_3} \right]$$

$$\nabla \times \mathbf{u} \equiv \text{curl } \mathbf{u} = \frac{\mathbf{i}_1}{h_2 h_3} \left[\frac{\partial(u_3 h_3)}{\partial q_2} - \frac{\partial(u_2 h_2)}{\partial q_3} \right] + \dots$$

$$\nabla \cdot \nabla\psi = \nabla^2\psi = \frac{1}{h_1 h_2 h_3} \left[\frac{\partial}{\partial q_1} \left(\frac{h_2 h_3}{h_1} \frac{\partial\psi}{\partial q_1} \right) + \dots \right]$$

$$\nabla(\nabla \cdot \mathbf{u}) = \frac{\mathbf{i}_1}{h_1} \frac{\partial}{\partial q_1} \left\{ \frac{1}{h_1 h_2 h_3} \left[\frac{\partial(u_1 h_2 h_3)}{\partial q_1} + \frac{\partial(u_2 h_3 h_1)}{\partial q_2} + \frac{\partial(u_3 h_1 h_2)}{\partial q_3} \right] \right\} + \dots$$

$$\nabla \times (\nabla \times \mathbf{u}) = \frac{\mathbf{i}_1}{h_2 h_3} \left\{ \frac{\partial}{\partial q_2} \left\{ \frac{h_3}{h_1 h_2} \left[\frac{\partial(u_2 h_2)}{\partial q_1} - \frac{\partial(u_1 h_1)}{\partial q_2} \right] \right\} - \frac{\partial}{\partial q_3} \left\{ \frac{h_2}{h_3 h_1} \left[\frac{\partial(u_1 h_1)}{\partial q_3} - \frac{\partial(u_3 h_3)}{\partial q_1} \right] \right\} \right\} + \dots$$

$$\begin{aligned} \nabla^2 \mathbf{u} &= \mathbf{i}_1 \left[\nabla^2 u_1 - u_1 h_1 \nabla^2 \frac{1}{h_1} \right. \\ &\quad + \frac{u_1 h_1}{h_1} \frac{\partial}{\partial q_1} (\nabla^2 q_1) + \frac{u_2 h_2}{h_1} \frac{\partial}{\partial q_1} (\nabla^2 q_2) + \frac{u_3 h_3}{h_1} \frac{\partial}{\partial q_1} (\nabla^2 q_3) \\ &\quad - \frac{2}{h_1^2} \frac{\partial(1/h_1)}{\partial q_1} \frac{\partial}{\partial q_1} (u_1 h_1) - \frac{2}{h_2^2} \frac{\partial(1/h_1)}{\partial q_2} \frac{\partial}{\partial q_1} (u_2 h_2) - \frac{2}{h_3^2} \frac{\partial(1/h_1)}{\partial q_3} \frac{\partial}{\partial q_1} (u_3 h_3) \\ &\quad \left. + \frac{1}{h_1} \frac{\partial(1/h_1^2)}{\partial q_1} \frac{\partial}{\partial q_1} (u_1 h_1) + \frac{1}{h_1} \frac{\partial(1/h_2^2)}{\partial q_1} \frac{\partial}{\partial q_2} (u_2 h_2) + \frac{1}{h_1} \frac{\partial(1/h_3^2)}{\partial q_1} \frac{\partial}{\partial q_3} (u_3 h_3) \right] + \dots \end{aligned}$$

$$\nabla \mathbf{u} = \sum_{j=1}^3 \sum_{k=1}^3 \frac{\mathbf{i}_j}{h_j} \mathbf{i}_k \left(\mathbf{i}_k \frac{\partial u_k}{\partial q_j} + u_k \frac{\partial \mathbf{i}_k}{\partial q_j} \right); \quad \frac{\partial \mathbf{i}_k}{\partial q_j} = \frac{\mathbf{i}_j}{h_k} \frac{\partial h_j}{\partial q_k} \quad (j \neq k), \quad \frac{\partial \mathbf{i}_j}{\partial q_j} = -\frac{\mathbf{i}_k}{h_k} \frac{\partial h_j}{\partial q_k} - \frac{\mathbf{i}_l}{h_l} \frac{\partial h_j}{\partial q_l}$$

$$\begin{aligned} &= \frac{\mathbf{i}_1 \mathbf{i}_1}{h_1} \left[\frac{\partial u_1}{\partial q_1} + \frac{u_2}{h_2} \frac{\partial h_1}{\partial q_2} + \frac{u_3}{h_3} \frac{\partial h_1}{\partial q_3} \right] + \frac{\mathbf{i}_1 \mathbf{i}_2}{h_1} \left[\frac{\partial u_2}{\partial q_1} - \frac{u_1}{h_2} \frac{\partial h_1}{\partial q_2} \right] + \frac{\mathbf{i}_1 \mathbf{i}_3}{h_1} \left[\frac{\partial u_3}{\partial q_1} - \frac{u_1}{h_3} \frac{\partial h_1}{\partial q_3} \right] \\ &\quad + \frac{\mathbf{i}_2 \mathbf{i}_1}{h_2} \left[\frac{\partial u_1}{\partial q_2} - \frac{u_2}{h_1} \frac{\partial h_2}{\partial q_1} \right] + \frac{\mathbf{i}_2 \mathbf{i}_2}{h_2} \left[\frac{\partial u_2}{\partial q_2} + \frac{u_3}{h_3} \frac{\partial h_2}{\partial q_3} + \frac{u_1}{h_1} \frac{\partial h_2}{\partial q_1} \right] + \frac{\mathbf{i}_2 \mathbf{i}_3}{h_2} \left[\frac{\partial u_3}{\partial q_2} - \frac{u_2}{h_3} \frac{\partial h_2}{\partial q_3} \right] \\ &\quad + \frac{\mathbf{i}_3 \mathbf{i}_1}{h_3} \left[\frac{\partial u_1}{\partial q_3} - \frac{u_3}{h_1} \frac{\partial h_3}{\partial q_1} \right] + \frac{\mathbf{i}_3 \mathbf{i}_2}{h_3} \left[\frac{\partial u_2}{\partial q_3} - \frac{u_3}{h_2} \frac{\partial h_3}{\partial q_2} \right] + \frac{\mathbf{i}_3 \mathbf{i}_3}{h_3} \left[\frac{\partial u_3}{\partial q_3} + \frac{u_1}{h_1} \frac{\partial h_3}{\partial q_1} + \frac{u_2}{h_2} \frac{\partial h_3}{\partial q_2} \right] \end{aligned}$$

$$\phi = \sum_{j=1}^3 \sum_{k=1}^3 \mathbf{i}_j \mathbf{i}_k \phi_{jk} \quad (\text{Dyadic})$$

$$\phi^T = \sum_{j=1}^3 \sum_{k=1}^3 \mathbf{i}_j \mathbf{i}_k \phi_{kj} \quad (\text{Dyadic transpose})$$

4 Cylindrical/Spherical Vector Identities

4.1 Cylindrical

$$\begin{aligned}
 \nabla\psi &= \mathbf{i}_r \frac{\partial\psi}{\partial r} + \mathbf{i}_\theta \frac{1}{r} \frac{\partial\psi}{\partial\theta} + \mathbf{i}_z \frac{\partial\psi}{\partial z} \\
 \nabla^2\psi &= \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial\psi}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2\psi}{\partial\theta^2} + \frac{\partial^2\psi}{\partial z^2} \\
 \nabla \cdot \mathbf{u} &= \frac{1}{r} \frac{\partial}{\partial r} (ru_r) + \frac{1}{r} \frac{\partial u_\theta}{\partial\theta} + \frac{\partial u_z}{\partial z} \\
 \nabla \times \mathbf{u} &= \mathbf{i}_r \left(\frac{1}{r} \frac{\partial u_z}{\partial\theta} - \frac{\partial u_\theta}{\partial z} \right) + \mathbf{i}_\theta \left(\frac{\partial u_r}{\partial z} - \frac{\partial u_z}{\partial r} \right) + \mathbf{i}_z \left[\frac{1}{r} \frac{\partial}{\partial r} (ru_\theta) - \frac{1}{r} \frac{\partial u_r}{\partial\theta} \right] \\
 \nabla^2 \mathbf{u} &= \mathbf{i}_r \left(\nabla^2 u_r - \frac{2}{r^2} \frac{\partial u_\theta}{\partial\theta} - \frac{u_r}{r^2} \right) + \mathbf{i}_\theta \left(\nabla^2 u_\theta + \frac{2}{r^2} \frac{\partial u_r}{\partial\theta} - \frac{u_\theta}{r^2} \right) + \mathbf{i}_z \nabla^2 u_z \\
 \nabla \mathbf{u} &= \mathbf{i}_r \mathbf{i}_r \frac{\partial u_r}{\partial r} + \mathbf{i}_r \mathbf{i}_\theta \frac{\partial u_\theta}{\partial r} + \mathbf{i}_r \mathbf{i}_z \frac{\partial u_z}{\partial r} \\
 &\quad + \mathbf{i}_\theta \mathbf{i}_r \frac{1}{r} \left(\frac{\partial u_r}{\partial\theta} - u_\theta \right) + \mathbf{i}_\theta \mathbf{i}_\theta \frac{1}{r} \left(\frac{\partial u_\theta}{\partial\theta} + u_r \right) + \mathbf{i}_\theta \mathbf{i}_z \frac{1}{r} \frac{\partial u_z}{\partial\theta} \\
 &\quad + \mathbf{i}_z \mathbf{i}_r \frac{\partial u_r}{\partial z} + \mathbf{i}_z \mathbf{i}_\theta \frac{\partial u_\theta}{\partial z} + \mathbf{i}_z \mathbf{i}_z \frac{\partial u_z}{\partial z}
 \end{aligned}$$

4.2 Spherical

$$\begin{aligned}
 \nabla\psi &= \mathbf{i}_R \frac{\partial\psi}{\partial R} + \mathbf{i}_\phi \frac{1}{R} \frac{\partial\psi}{\partial\phi} + \mathbf{i}_\theta \frac{1}{R \sin\phi} \frac{\partial\psi}{\partial\theta} \\
 \nabla^2\psi &= \frac{1}{R^2} \left[\frac{\partial}{\partial R} \left(R^2 \frac{\partial\psi}{\partial R} \right) + \frac{1}{\sin\phi} \frac{\partial}{\partial\phi} \left(\sin\phi \frac{\partial\psi}{\partial\phi} \right) + \frac{1}{\sin^2\phi} \frac{\partial^2\psi}{\partial\theta^2} \right] \\
 \nabla \cdot \mathbf{u} &= \frac{1}{R^2} \frac{\partial}{\partial R} (R^2 u_R) + \frac{1}{R \sin\phi} \frac{\partial}{\partial\phi} (\sin\phi u_\phi) + \frac{1}{R \sin\phi} \frac{\partial u_\theta}{\partial\theta} \\
 \nabla \times \mathbf{u} &= \mathbf{i}_R \frac{1}{R \sin\phi} \left[\frac{\partial}{\partial\phi} (\sin\phi u_\theta) - \frac{\partial u_\phi}{\partial\theta} \right] + \mathbf{i}_\phi \frac{1}{R} \left[\frac{1}{\sin\phi} \frac{\partial u_R}{\partial\theta} - \frac{\partial}{\partial R} (R u_\theta) \right] + \mathbf{i}_\theta \frac{1}{R} \left[\frac{\partial}{\partial R} (R u_\phi) - \frac{\partial u_R}{\partial\phi} \right] \\
 \nabla^2 \mathbf{u} &= \mathbf{i}_R \left[\nabla^2 u_R - \frac{2}{R^2} \left(u_R + \frac{\partial u_\phi}{\partial\phi} + u_\phi \cot\phi - \frac{1}{\sin\phi} \frac{\partial u_\theta}{\partial\theta} \right) \right] \\
 &\quad + \mathbf{i}_\phi \left[\nabla^2 u_\phi + \frac{1}{R^2} \left(2 \frac{\partial u_R}{\partial\phi} - \frac{u_\phi}{\sin^2\phi} - 2 \frac{\cos\phi}{\sin^2\phi} \frac{\partial u_\theta}{\partial\theta} \right) \right] \\
 &\quad + \mathbf{i}_\theta \left[\nabla^2 u_\theta + \frac{1}{R^2} \left(\frac{2}{\sin\phi} \frac{\partial u_R}{\partial\theta} + 2 \frac{\cos\phi}{\sin^2\phi} \frac{\partial u_\phi}{\partial\theta} - \frac{u_\theta}{\sin^2\phi} \right) \right] \\
 \nabla \mathbf{u} &= \mathbf{i}_R \mathbf{i}_R \frac{\partial u_R}{\partial R} + \mathbf{i}_R \mathbf{i}_\phi \frac{\partial u_\phi}{\partial R} + \mathbf{i}_R \mathbf{i}_\theta \frac{\partial u_\theta}{\partial R} \\
 &\quad + \mathbf{i}_\phi \mathbf{i}_R \frac{1}{R} \left(\frac{\partial u_R}{\partial\phi} - u_\phi \right) + \mathbf{i}_\phi \mathbf{i}_\phi \frac{1}{R} \left(\frac{\partial u_\phi}{\partial\phi} + u_R \right) + \mathbf{i}_\phi \mathbf{i}_\theta \frac{1}{R} \frac{\partial u_\theta}{\partial\phi} \\
 &\quad + \mathbf{i}_\theta \mathbf{i}_R \frac{1}{R} \left(\frac{1}{\sin\phi} \frac{\partial u_R}{\partial\theta} - u_\theta \right) + \mathbf{i}_\theta \mathbf{i}_\phi \frac{1}{R} \left(\frac{1}{\sin\phi} \frac{\partial u_\phi}{\partial\theta} - \cot\phi u_\theta \right) + \mathbf{i}_\theta \mathbf{i}_\theta \frac{1}{R} \left(\frac{1}{\sin\phi} \frac{\partial u_\theta}{\partial\theta} + u_R + \cot\phi u_\phi \right)
 \end{aligned}$$

5 Vector Integral Theorems

5.1 Gauss's Divergence Theorem

$$\oiint \mathbf{f} \cdot d\mathbf{S} = \iiint \nabla \cdot \mathbf{f} dV$$

5.2 Green's Theorems

5.2.1 Green's First Theorem

$$\oiint \phi \frac{\partial \psi}{\partial n} dS = \iiint (\phi \nabla^2 \psi + \nabla \phi \cdot \nabla \psi) dV$$

5.2.2 Green's Second Theorem

$$\oiint \left(\phi \frac{\partial \psi}{\partial n} - \psi \frac{\partial \phi}{\partial n} \right) dS = \iiint (\phi \nabla^2 \psi - \psi \nabla^2 \phi) dV$$

5.2.3 Special Cases

$$\begin{aligned} \oiint (\phi \nabla \phi) \cdot d\mathbf{S} &= \iiint [\phi \nabla^2 \phi + (\nabla \phi)^2] dV \\ \oiint \frac{\partial \phi}{\partial n} dS &= \iiint \nabla^2 \phi dV \end{aligned}$$

5.3 Stokes' Theorem

Let a simple closed curve C be spanned by a surface S . Define the positive normal \mathbf{n} to S , and the positive sense of description of the curve C with line element $d\mathbf{r}$, such that the positive sense of the contour C is clockwise when we look through the surface S in the direction of the normal. Then, if \mathbf{f} is continuously differentiable vector field defined on S and C with vector element $\mathbf{S} = \mathbf{n}dS$

$$\oint \mathbf{f} \cdot d\mathbf{r} = \iint \nabla \times \mathbf{f} \cdot d\mathbf{S}$$

where the line integral around C is taken in the positive sense.

5.4 Integral rate of change theorems

5.4.1 Rate of change of volume integral bounded by a moving closed surface.

Let f be a continuous scalar function of position and time t defined throughout the volume $V(t)$, which is itself bounded by a simple closed surface $S(t)$ moving with velocity \mathbf{v} . Then the rate of change of the volume integral of f is given by

$$\frac{D}{Dt} \int_{V(t)} f dV = \int_{V(t)} \frac{\partial f}{\partial t} dV + \int_{S(t)} f \mathbf{v} \cdot d\mathbf{S}$$

where $d\mathbf{S}$ is the outward drawn vector element of area, and

$$\frac{D}{Dt} \equiv \frac{\partial}{\partial t} + (\mathbf{v} \cdot \nabla)$$

5.4.2 Rate of change of flux through a surface.

Let \mathbf{q} be a vector function that may also depend on the time t , and \mathbf{n} be the unit outward drawn normal to the surface S that moves with velocity \mathbf{v} . Defining the flux of \mathbf{q} through S as

$$m = \oint_S \mathbf{q} \cdot \mathbf{n} dS$$

then

$$\frac{Dm}{Dt} = \oint \left[\frac{\partial \mathbf{q}}{\partial t} + \mathbf{v}(\nabla \cdot \mathbf{q}) + \nabla \times (\mathbf{q} \times \mathbf{v}) \right] \cdot \mathbf{n} dS$$

5.4.3 Rate of change of the circulation around a given moving curve.

Let C be a closed curve, moving with velocity \mathbf{v} , on which is defined a vector field \mathbf{q} . Defining the circulation ζ of \mathbf{q} around C by

$$\zeta = \oint \mathbf{q} \cdot d\mathbf{r}$$

then

$$\frac{D\zeta}{Dt} = \oint \left[\frac{\partial \mathbf{q}}{\partial t} + (\nabla \times \mathbf{q}) \times \mathbf{v} \right] \cdot d\mathbf{r}$$

6 Ideal Flow (irrotational/inviscid/incompressible)

6.1 Governing equations

$$\begin{aligned}\nabla \cdot \mathbf{u} &= 0 && \text{(continuity)} \\ \rho \left[\frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} \right] &= -\nabla p && \text{(momentum)}\end{aligned}$$

6.2 Fundamentals

6.2.1 Streamfunction

May be defined when continuity eq. reduces to two terms.
For cartesian coordinates:

$$\frac{\partial u}{\partial x} = -\frac{\partial v}{\partial y}$$

The streamfunction may be defined as:

$$u \equiv \frac{\partial \psi}{\partial y} \quad \text{and} \quad v \equiv -\frac{\partial \psi}{\partial x}$$

6.2.2 Velocity potential

$$u \equiv \frac{\partial \phi}{\partial x} \quad \text{and} \quad v \equiv \frac{\partial \phi}{\partial y}$$

6.2.3 Bernoulli's equation

If flow is irrotational throughout:

$$p + \frac{1}{2}\rho V^2 = \text{const.}$$

else, only valid along streamlines.

6.2.4 Incompressible pressure coefficient

$$C_p = 1 - \left(\frac{V}{V_\infty} \right)^2$$

7 Viscous Flow

7.1 Governing Equations

$$\nabla \cdot \mathbf{u} = 0 \quad (\text{continuity})$$

$$\rho(\mathbf{u} \cdot \nabla \mathbf{u}) = \rho \mathbf{F} - \nabla p + \mu \nabla^2 \mathbf{u} \quad (\text{momentum})$$

7.2 Boundary layer terms

7.2.1 Reynolds Numbers

$$\text{Re} = \frac{\rho U l}{\mu}$$

7.2.2 Skin friction coefficient

$$C_f(x) = \frac{\tau_w(x)}{\frac{1}{2} \rho U^2}$$

$$= \frac{2}{\sqrt{\text{Re}}} \left(\frac{\partial u^*}{\partial y} \right)_w$$

7.2.3 Drag coefficient

$$C_D = \frac{F_D}{\frac{1}{2} \rho U^2 A}$$

$$= b \int_{\text{LE}}^{\text{TE}} (-p_u \sin \theta + \tau_u \cos \theta) dS_u$$

$$+ b \int_{\text{LE}}^{\text{TE}} (p_l \sin \theta + \tau_l \cos \theta) dS_l$$

where "u", "l", and "b"
denote "upper", "lower", and "width"

7.2.4 Separation point

$$\left(\frac{\partial u}{\partial y} \right)_{y=0} = 0$$

7.2.5 Displacement thickness

$$\delta^* = \int_{y=0}^{\infty} \left(1 - \frac{u}{U_e} \right) dy$$

7.2.6 Momentum thickness

$$\theta = \int_{y=0}^{\infty} \frac{u}{U_e} \left(1 - \frac{u}{U_e} \right) dy$$

7.2.7 Shape factor

$$H = \frac{\delta^*}{\theta}$$

7.3 Blasius equation

Assumptions:

- Incompressible
- Laminar
- Steady
- $\frac{\partial p}{\partial x} = 0$

7.3.1 Equation

$$f''' + \frac{1}{2} f f'' = 0 \quad \text{where} \quad \eta = \sqrt{vx}/U$$

7.3.2 Shear stress

$$\tau_w = \mu \left(\frac{du}{dy} \right)_{y=0} = 0.332 \rho U^2 / \sqrt{\text{Re}_x}$$

7.3.3 Skin friction coefficient

$$C_f \equiv \frac{\tau_w}{\frac{1}{2} \rho U^2} = \frac{0.664}{\sqrt{\text{Re}_x}}$$

7.3.4 Drag on plate

$$F_D = \int_0^L \tau_w dx = \frac{0.664 \rho U^2 L}{\sqrt{\text{Re}_L}}$$

(drag force per unit width)

7.3.5 Drag coefficient

$$C_D = \frac{F_D}{\frac{1}{2} \rho U^2 L} = \frac{1.33}{\sqrt{\text{Re}_L}}$$

7.4 Pohlhausen's Polynomial

$$\frac{u(\eta)}{U_\infty} = \left(2 + \frac{1}{6} \Lambda \right) \eta - \frac{1}{2} \Lambda \eta^2 + \left(\frac{1}{2} \Lambda - 2 \right) \eta^3 + \left(1 - \frac{1}{6} \Lambda \right) \eta^4$$

where: $\eta = y/\delta \quad \Lambda = \frac{\rho \delta^2}{\mu} \frac{dU}{dx}$

7.5 Von Karman momentum integral equation

$$\tau_w = \rho \left[\frac{d}{dx} (U_e^2 \theta) + U_e \delta^* \frac{dU_e}{dx} \right]$$

$$\tau_w = \mu \left. \frac{\partial u}{\partial y} \right|_{y \rightarrow 0}$$

solve for $\delta(x)$ (boundary layer thickness)

$$C_f = \frac{\tau_w}{\frac{1}{2} \rho U^2} \quad (\text{skin friction coefficient})$$

$$\text{Re} = \frac{\rho U \ell}{\mu} \quad \text{and} \quad \text{Re}_x = \frac{\rho U x}{\mu}$$

(Reynolds number and x-Reynolds number)

7.6 Thwaites' method

For solving Von Karman mom. int. eq.:

$$\theta^2 = \frac{0.45\nu}{U_e^2(x)} \int_0^x U_e^5(x') dx' + \frac{\theta_0^2 U_0^6}{U_e^6(x)}$$

$$\lambda \equiv \frac{\theta^2}{\nu} \frac{dU_e}{dx} \quad (\text{solve for } \theta)$$

$$\tau \equiv \mu \frac{U_e}{\theta} l(\lambda)$$

$$H(\lambda) = \frac{\delta^*}{\theta}$$

(use table to find surface shear stress and displacement thickness)

8 Compressible Flow

8.1 Governing equations

$$\frac{\partial}{\partial t} \iiint_V \rho dV + \iint_S \rho \mathbf{u} \cdot d\mathbf{S} = 0 \quad (\text{continuity})$$

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{u}) = 0$$

$$\frac{\partial}{\partial t} \iiint_V \rho \mathbf{u} dV + \iint_S \rho (\mathbf{u} \cdot d\mathbf{S}) \mathbf{u} = \iiint_V \rho \mathbf{F} dV - \iint_S p d\mathbf{S} \quad (\text{momentum})$$

$$\rho \left(\frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} \right) = \rho \mathbf{F} - \nabla p$$

8.2 General Definitions

8.2.1 Mach number

$$M \equiv U/c$$

where $c^2 \equiv \left(\frac{\partial p}{\partial \rho} \right)_s$

8.2.2 Compressibility

$$\tau = -\frac{1}{V} \frac{dV}{dp}$$

$$= \frac{1}{\rho} \frac{d\rho}{dp}$$

8.2.3 Isothermal compressibility

$$\tau_T = -\frac{1}{V} \left(\frac{\partial V}{\partial p} \right)_T$$

8.2.4 Isentropic compressibility

$$\tau_s = -\frac{1}{V} \left(\frac{\partial V}{\partial p} \right)_s$$

8.3 Perfect Gas Thermodynamic Relations

8.3.1 Internal Energy/Enthalpy/Thermal eq. of state

$$e = c_v T \quad (\text{internal energy})$$

$$h = c_p T \quad (\text{enthalpy})$$

$$p = \rho RT \quad (\text{thermal equation of state})$$

8.3.2 Specific heats

$$c_v = \frac{R}{\gamma - 1} \quad c_p = \frac{\gamma R}{\gamma - 1} \quad \gamma = c_p / c_v$$

8.3.3 Speed of sound

$$c = \sqrt{\gamma RT} = \sqrt{\gamma p / \rho}$$

8.3.4 Entropy change

$$s_2 - s_1 = c_p \ln \left(\frac{T_2}{T_1} \right) - R \ln \left(\frac{p_2}{p_1} \right)$$

$$= c_v \ln \left(\frac{T_2}{T_1} \right) - R \ln \left(\frac{\rho_2}{\rho_1} \right)$$

8.3.5 Isentropic relations

$$\frac{p_2}{p_1} = \left(\frac{\rho_2}{\rho_1} \right)^\gamma$$

$$\frac{T_2}{T_1} = \left(\frac{\rho_2}{\rho_1} \right)^{\gamma-1} = \left(\frac{p_2}{p_1} \right)^{\frac{\gamma-1}{\gamma}}$$

8.3.6 Important air properties

$$R = 287 \text{ m}^2 / (\text{s}^2 \text{K})$$

$$c_v = 717 \text{ m}^2 / (\text{s}^2 \text{K})$$

$$c_p = 1004 \text{ m}^2 / (\text{s}^2 \text{K})$$

$$\gamma = 1.40$$

8.4 One-dimensional isentropic flow

8.4.1 Continuity

$$\rho_1 u_1 A_1 = \rho_2 u_2 A_2$$

8.4.2 Stagnation enthalpy

$$h_0 = h + \frac{1}{2} u^2 = \text{const.}$$

8.4.3 Stagnation property relations

$$\begin{aligned}\frac{T_0}{T} &= 1 + \frac{\gamma-1}{2}M^2 \\ \frac{p_0}{p} &= \left(1 + \frac{\gamma-1}{2}M^2\right)^{\frac{\gamma}{\gamma-1}} \\ \frac{\rho_0}{\rho} &= \left(1 + \frac{\gamma-1}{2}M^2\right)^{\frac{1}{\gamma-1}}\end{aligned}$$

8.4.4 Nozzle area relation

$$\frac{A}{A^*} = \frac{1}{M} \left[\frac{2}{\gamma+1} \left(1 + \frac{\gamma-1}{2}M^2\right) \right]^{\frac{\gamma+1}{2(\gamma-1)}}$$

8.4.5 Velocity-area differential relationship

$$\frac{du}{u} = -\frac{1}{1-M^2} \frac{dA}{A}$$

8.5 Normal shock waves

8.5.1 Mach number, pressure, temperature, and density relations across a shock

$$\begin{aligned}M_2^2 &= \frac{(\gamma-1)M_1^2 + 2}{2\gamma M_1^2 + 1 - \gamma} \\ \frac{p_2}{p_1} &= 1 + \frac{2\gamma}{\gamma+1}(M_1^2 - 1) \\ \frac{\rho_2}{\rho_1} &= \frac{(\gamma+1)M_1^2}{(\gamma-1)M_1^2 + 2} \\ \frac{T_2}{T_1} &= 1 + \frac{2(\gamma-1)}{(\gamma+1)^2} \frac{\gamma M_1^2 + 1}{M_1^2} (M_1^2 - 1)\end{aligned}$$

8.5.2 Entropy change

$$\begin{aligned}\frac{s_2 - s_1}{c_v} &= \ln \left\{ \frac{p_2}{p_1} \left(\frac{\rho_1}{\rho_2} \right)^\gamma \right\} \\ &= \ln \left\{ \left[1 + \frac{2\gamma}{\gamma-1} (M_1^2 - 1) \right] \left[\frac{(\gamma-1)M_2^2 + 2}{(\gamma+1)M_1^2} \right]^\gamma \right\}\end{aligned}$$

8.6 Wind Tunnel equations

8.6.1 Mass flow rate

$$\dot{m} = \frac{\rho_0 A^*}{T_0} \sqrt{\frac{\gamma}{R} \left(\frac{2}{\gamma+1} \right)^{\frac{\gamma+1}{\gamma-1}}}$$

8.6.2 Nozzle/Diffuser area ratio

$$\begin{aligned}\frac{A_2^*}{A_1^*} &= \frac{p_0}{p'_0} \\ &= \left[1 + \frac{2\gamma}{\gamma+1} (M_1^2 - 1) \right]^{\frac{1}{\gamma-1}} \left[\frac{(\gamma-1)M_1^2 + 2}{(\gamma+1)M_1^2} \right]^{\frac{\gamma}{\gamma-1}}\end{aligned}$$

where A_2^* , A_1^* , p_0 , p'_0 , and M_1 are the diffuser throat area, inflow nozzle throat area, stagnation pressure behind the shock, stagnation pressure after the shock, and Mach number upstream of the shock

8.7 Two-dimensional compressible flow

8.7.1 Definitions

δ = deflection angle

σ = wave angle

M_1 = upstream Mach number

M_2 = downstream Mach number

8.7.2 Mach wave angle

$$\sin \mu = 1/M$$

8.7.3 Mach number relation

$$M_2^2 \sin^2(\sigma - \delta) = \frac{(\gamma-1)M_1^2 \sin^2 \sigma + 2}{2\gamma M_1^2 \sin^2 \sigma + 1 - \gamma}$$

8.7.4 Oblique shockwave angle

$$\tan \delta = 2 \cot \sigma \frac{M_1^2 \sin^2 \sigma - 1}{M_1^2 (\gamma - \cos 2\sigma) + 2}$$

8.7.5 Prandtl-Meyer function (expansion fans)

$$\begin{aligned}\delta_1 + \nu(M_1) &= \delta_2 + \nu(M_2) = \text{const.} \\ \nu(M) &= \sqrt{\frac{\gamma+1}{\gamma-1}} \arctan \sqrt{\frac{\gamma-1}{\gamma+1} (M^2 - 1)} - \arctan \sqrt{M^2 - 1}\end{aligned}$$

8.8 Thin-airfoil theory

8.8.1 Lift and Drag coefficient

$$C_L = \frac{L}{\frac{1}{2}\rho_\infty U_\infty^2 b}$$
$$\approx \frac{4\alpha}{\sqrt{M_\infty^2 - 1}}$$
$$C_D = \frac{D}{\frac{1}{2}\rho_\infty U_\infty^2 b}$$
$$\approx \frac{4\alpha^2}{\sqrt{M_\infty^2 - 1}}$$

8.9 Velocity Potential

$$\nabla^2 \phi = \frac{1}{2}M_0^2 [(\gamma - 1)\nabla^2 \phi (\nabla \phi \cdot \nabla \phi) + \nabla \phi \cdot \nabla (\nabla \phi \cdot \nabla \phi)]$$

where

$$a^2 = a_0^2 - \frac{\gamma - 1}{2}V^2 = a_0^2 - \frac{\gamma - 1}{2}(u^2 + v^2 + w^2)$$

9 Rocket Propulsion Equations

9.1 Definitions and Fundamentals

9.1.1 Total impulse

$$I_t = \int_0^t F dt$$

$$= Ft \quad (\text{for constant thrust})$$

9.1.2 Specific impulse

$$I_s = I_t / (m_p g_0)$$

$$= F / (\dot{m} g_0)$$

$$= F / \dot{w}$$

$$= I_t / w$$

9.1.3 Effective exhaust velocity

$$c = I_s g_0 = F / \dot{m}$$

9.1.4 Mass ratio MR

$$m_f / m_0$$

9.1.5 Propellant mass fraction ζ

$$\zeta = m_p / m_0$$

$$= (m_0 - m_f) / m_0 = m_p / (m_p + m_f)$$

where $m_0 = m_p + m_f$

9.1.6 Impulse-to-weight Ratio

$$\frac{I_t}{w_0} = \frac{I_t}{(m_f + m_p) g_0} = \frac{I_s}{m_f / m_p + 1}$$

9.1.7 Thrust F

$$F = \frac{d(mv_2)}{dt} = \dot{m}v_2 \quad \text{at sea level} = \frac{\dot{w}}{g_0} v_2$$

$$= \dot{m}v_2 + (p_2 - p_3) A_2$$

$$= \dot{m}c$$

$$= \dot{m}v_2 + p_2 A_2 \quad (\text{in vacuum of space})$$

$$= F_{\text{opt}} + p_1 A_t \left(\frac{p_2}{p_1} - \frac{p_3}{p_1} \right) \epsilon$$

$$= \frac{A_t v_t v_2}{V_t} + (p_2 - p_3) A_2$$

$$= A_t p_1 \sqrt{\frac{2\gamma^2}{\gamma-1} \left(\frac{2}{\gamma+1} \right)^{\frac{\gamma+1}{\gamma-1}} \left[1 - \left(\frac{p_2}{p_1} \right)^{\frac{\gamma-1}{\gamma}} \right]} + (p_2 - p_3) A_2$$

$$= C_F A_t p_1$$

$$= \dot{m} c^* C_F$$

9.1.8 Thrust Coefficient

$$C_F = \frac{F}{p_1 A_t}$$

$$= \sqrt{\frac{2\gamma^2}{\gamma-1} \left(\frac{2}{\gamma+1} \right)^{\frac{\gamma+1}{\gamma-1}} \left[1 - \left(\frac{p_2}{p_1} \right)^{\frac{\gamma-1}{\gamma}} \right]} + \frac{p_2 - p_3}{p_1} \frac{A_2}{A_t}$$

9.1.9 Exhaust velocity c

$$c = v_2 + (p_2 - p_3) A_2 / \dot{m} = I_s g_0$$

9.1.10 Characteristic velocity c^*

$$c^* = p_1 A_t / \dot{m}$$

$$= \frac{I_s g_0}{C_F}$$

$$= \frac{c}{C_F}$$

$$= \frac{1}{\gamma} \sqrt{\frac{\gamma R T_1}{\left(\frac{2}{\gamma+1} \right)^{\frac{\gamma+1}{\gamma-1}}}}$$

9.1.11 Power of the jet P_{jet}

$$P_{\text{jet}} = \frac{1}{2} \dot{m} v_2^2 = \frac{1}{2} \dot{w} g_0 I_s^2 = \frac{1}{2} F g_0 I_s = \frac{1}{2} F v_2$$

9.1.12 Power input to a chemical engine

$$P_{\text{chem}} = \dot{m} Q_R$$

where Q_R is the heat of combustion per unit of propellant mass.

9.1.13 Power transmitted to vehicle

$$P_{\text{vehicle}} = Fu \quad (\text{where } u \text{ is vehicle velocity})$$

9.1.14 Internal efficiency η_p

$$\eta_{\text{int}} = \frac{\text{kinetic power in jet}}{\text{available chemical power}} = \frac{\frac{1}{2} \dot{m} v_2^2}{\eta_{\text{comp}} P_{\text{chem}}}$$

9.1.15 Propulsive efficiency

$$\eta_p = \frac{\text{vehicle power}}{\text{vehicle power} + \text{residual kinetic jet power}}$$

$$= \frac{Fu}{Fu + \frac{1}{2} \dot{m} (c + u)^2} = \frac{2u/c}{1 + (u/c)^2}$$

9.1.16 Multiple propulsion systems

$$F_{\text{oa}} = \sum F = F_1 + F_2 + F_3 + \dots$$

$$\dot{m}_{\text{oa}} = \sum \dot{m} = \dot{m}_1 + \dot{m}_2 + \dot{m}_3 + \dots$$

$$(I_s)_{\text{oa}} = F_{\text{oa}} / (g_0 \dot{m}_{\text{oa}})$$

9.2 Nozzle theory and thermodynamic relations

9.2.1 Stagnation enthalpy

$$h_0 = h + \frac{1}{2}v^2 = \text{const.}$$

9.2.2 Perfect gas law

$$p_x V_x = RT_x$$

9.2.3 Specific heats c and their γ

$$k = c_p / c_v$$

$$c_p - c_v = R$$

$$c_p = \gamma R / (\gamma - 1)$$

9.2.4 Isentropic flow between x and y in nozzle

$$T_x / T_y = (p_x / p_y)^{\frac{\gamma-1}{\gamma}} = (V_y / V_x)^{\gamma-1}$$

9.2.5 Stagnation temperature

$$T_0 = T + \frac{1}{2}v^2 / c_p$$

9.2.6 Speed of sound & Mach number

$$a = \sqrt{\gamma RT}$$

$$M = v/a = v / \sqrt{\gamma RT}$$

9.2.7 Isentropic relations (temperature, Mach number, pressure)

$$T_0 = T \left[1 + \frac{1}{2}(\gamma - 1)M^2 \right]$$

$$M = \sqrt{\frac{2}{\gamma - 1} \left(\frac{T_0}{T} - 1 \right)}$$

$$p_0 = p \left[1 + \frac{1}{2}(\gamma - 1)M^2 \right]^{\frac{\gamma}{\gamma - 1}}$$

9.2.8 General area ratio

$$\frac{A_y}{A_x} = \frac{M_x}{M_y} \sqrt{\frac{1 + M_y^2(\gamma - 1)/2}{1 + M_x^2(\gamma - 1)/2}}$$

9.2.9 Exit velocity

$$v_2 = \sqrt{2(h_1 - h_2) + v_1^2}$$

$$= \sqrt{\frac{2\gamma}{\gamma - 1} RT_1 \left[1 - \left(\frac{p_2}{p_1} \right)^{(\gamma-1)/\gamma} \right] + v_1^2}$$

$$= \sqrt{\frac{2\gamma}{\gamma - 1} RT_1 \left[1 - \left(\frac{p_2}{p_1} \right)^{(\gamma-1)/\gamma} \right]} \quad (\text{when } v_1 \approx 0(\text{chamber}))$$

$$= \sqrt{\frac{2\gamma}{\gamma - 1} \frac{R'T_1}{M} \left[1 - \left(\frac{p_2}{p_1} \right)^{(\gamma-1)/\gamma} \right]}$$

where M is molecular mass.

9.2.10 Critical/Throat velocity

$$v_t = \sqrt{\frac{2\gamma}{\gamma + 1} RT_1} = \sqrt{\gamma RT_t} = a_t$$

9.2.11 Mass flow rate

$$\dot{m} = \frac{A_t v_t}{V_t} = A_t p_1 \gamma \sqrt{\frac{\left(\frac{2}{\gamma+1} \right)^{\frac{\gamma+1}{\gamma-1}}}{\gamma RT_1}}$$

9.2.12 Ratio between throat and downstream area with pressure p_y

$$\frac{A_t}{A_y} = \frac{V_t v_y}{V_y v_t} = \left(\frac{\gamma + 1}{2} \right)^{\frac{1}{\gamma-1}} \left(\frac{p_y}{p_1} \right)^{1/\gamma} \sqrt{\frac{\gamma + 1}{\gamma - 1} \left[1 - \left(\frac{p_y}{p_1} \right)^{\frac{\gamma-1}{\gamma}} \right]}$$

9.2.13 Velocity of downstream point from throat

$$\frac{v_y}{v_t} = \sqrt{\frac{\gamma + 1}{\gamma - 1}}$$

9.2.14 Cone-shaped nozzle correction factor for exhaust velocity

$$\gamma = \frac{1}{2}(1 + \cos \alpha)$$

where α is the cone angle.

9.2.15 Rocket Equation

$$e^{\Delta u/c} = 1/\mathbf{MR} = m_0/m_f$$

9.3 Chemical Rocket Performance Analysis

9.3.1 Mol fraction

$$X_j = \frac{n_j}{n} \quad \text{and} \quad n = \sum_{j=1}^m n_j$$

9.3.2 Effective average molecular mass

$$\mathbb{M} = \frac{\sum_{j=1}^m n_j \mathbb{M}_j}{\sum_{j=1}^m n_j}$$

9.3.3 Molar specific heat and specific heat ratio

$$(C_p)_{\text{mix}} = \frac{\sum_{j=1}^m n_j (C_p)_j}{\sum_{j=1}^m n_j}$$

$$\gamma_{\text{mix}} = \frac{(C_p)_{\text{mix}}}{(C_p)_{\text{mix}} - R'}$$

9.3.4 Heat of reaction

$$\Delta_r H^0 = \sum [n_j (\Delta_f H^0)]_{\text{products}} - \sum [n_j (\Delta_f H^0)]_{\text{reactants}}$$

9.3.5 Gibbs free energy

$$G = \sum_{j=1}^m n_j G_j$$

$$G_j = U_j + p_j V_j - T_j S_j = h_j - T_j S_j$$

$$\Delta_r G^0 = \sum_{j=1}^m [n_j (\Delta_f G^0)]_{\text{products}} - \sum_{j=1}^r [n_j (\Delta_f G^0)]_{\text{reactants}}$$

9.3.6 Heat of reaction

$$\Delta_r H = \sum_1^m n_j \int_{T_{\text{ref}}}^{T_1} C_{p_j} dT = \sum_1^m n_j \Delta h_j \Big|_{T_{\text{ref}}}^{T_1}$$

9.4 Solid Propellant Rocket Motor Fundamentals

9.4.1 Mass flow rate

$$\begin{aligned} \dot{m} &= A_b r \rho_b \\ &= \frac{d(\rho_1 V_1)}{dt} + \frac{A_t p_1}{c^*} \\ &= A_b \rho_b a p_1^n \end{aligned}$$

where A_b , r , and ρ_b are the burn area, burn rate, and density of solid propellant before burning

9.4.2 Pressure in steady burning

$$p_1 = K \rho_b r c^* \quad \text{where} \quad K \equiv A_b / A_t$$

9.4.3 Burning rate approximation

$$r = a p_1^n$$

10 Aeroacoustics

10.1 Governing Equations

$$\frac{\partial p'}{\partial t} + \rho_0 \nabla \cdot \mathbf{u}' = 0 \quad (\text{continuity})$$

$$\rho_0 \frac{\partial \mathbf{u}'}{\partial t} = -\nabla p' \quad (\text{momentum})$$

10.2 General definitions

10.2.1 Frequency

$$f = \frac{\omega}{2\pi} = \frac{1}{T} \rightarrow \omega = \frac{1}{T}$$

f = frequency

ω = angular rate

T = period

10.2.2 Fourier Series

$$f(t) = \frac{a_0}{2} + \sum_{n=1}^{\infty} \left[a_n \cos\left(\frac{2\pi n t}{T}\right) + b_n \sin\left(\frac{2\pi n t}{T}\right) \right]$$

$$a_n = \frac{2}{T} \int_0^T f(t) \cos\left(\frac{2\pi n t}{T}\right) dt$$

$$b_n = \frac{2}{T} \int_0^T f(t) \sin\left(\frac{2\pi n t}{T}\right) dt$$

$$f(t) = \sum_{n=-\infty}^{\infty} c_n e^{i \frac{2\pi n t}{T}}$$

$$c_n = \frac{1}{T} \int_0^T f(t) e^{i \frac{2\pi n t}{T}} dt$$

$$c_n = \begin{cases} \frac{a_{-n} + i b_{-n}}{2}; & n \leq -1 \\ \frac{a_0}{2} & n = 0 \\ \frac{a_n - i b_n}{2} & n \geq 1 \end{cases}$$

10.2.3 Fourier Transform

$$\bar{f}(\omega) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(t) e^{-i\omega t} dt \quad (\text{fourier transform})$$

$$f(t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \bar{f}(\omega) e^{i\omega t} d\omega \quad (\text{inverse fourier transform})$$

10.2.4 Acoustic Intensity

$$\mathbf{I} = p' \mathbf{u}'$$

$$\langle I \rangle = \frac{A^2}{2\rho_0 a_0} \quad \text{Units: } \left[\frac{\text{W}}{\text{m}^2} \right]$$

10.2.5 Acoustic Energy

$$e = \frac{p'^2}{2\rho_0 a_0^2} + \frac{1}{2} \rho_0 \mathbf{u}'^2$$

$$\langle e \rangle = \frac{A^2}{2\rho_0 a_0^2} \quad \text{Units: } \left[\frac{\text{J}}{\text{m}^3} \right] \quad (\text{energy density})$$

10.2.6 Root of Square Average (RMS)

$$p'_{\text{rms}} = \sqrt{\frac{1}{T} \int_0^T p(t)^2 dt}$$

$$= \frac{A}{\sqrt{2}} \quad (\text{for } p(t)' = A \cos t)$$

10.2.7 Decibel

$$10 \log_{10} \left(\frac{\text{Power}}{\text{Ref power}} \right)$$

10.2.8 Acoustic Power Level

$$\text{APL} = 10 \log_{10} \left(\frac{\mathbb{P}_{\text{acoustic}}}{\mathbb{P}_{\text{ref}}} \right) \quad [\text{dB}]$$

$$\mathbb{P} = \mathbb{P}_{\text{ref}} 10^{\text{APL}/10}$$

where $\mathbb{P}_{\text{ref}} = 10^{-12}$ [W] and $\mathbb{P} = \langle I \rangle A$

where A is the area the sound is going through

10.2.9 Sound Pressure Level

$$\text{SPL} = 10 \log_{10} \left(\frac{p'_{\text{rms}}{}^2}{p_{\text{ref}}^2} \right)$$

$$= 20 \log_{10} \left(\frac{p_{\text{rms}}}{p_{\text{ref}}} \right)$$

$$p_{\text{ref}} = 2 \times 10^{-5} \text{ [Pa]}$$

10.2.10 RMS pressure summation

$$(p_{\text{rms}})_{1+2} = \sqrt{(p_{\text{rms}})_1^2 + (p_{\text{rms}})_2^2} \quad (1)$$

10.2.11 SPL summation

$$\text{SPL}_{1+2+\dots+n} = 10 \log_{10} \left(10^{\text{SPL}_1/10} + 10^{\text{SPL}_2/10} + \dots + 10^{\text{SPL}_n/10} \right) \quad (2)$$

10.2.12 Impedence

$$Z = i\sigma\omega$$

$$= \frac{p'}{u'}$$

$$m/A = \sigma \quad (\text{mass/area})$$

10.2.13 Planar wave transmission & reflection

$$p_t = \frac{2\rho_0 a_0}{2\rho_0 a_0 + i\sigma\omega} p_i \quad (\text{transmission})$$

$$p_r = \frac{i\sigma\omega}{2\rho_0 a_0 + i\sigma\omega} p_i \quad (\text{reflection})$$

$$\text{where } \sigma = m/A$$

10.2.14 Reflection/Transmission coefficient

$$\alpha_r \equiv \frac{\langle I \rangle_r}{\langle I \rangle_i}$$

$$\alpha_t \equiv \frac{\langle I \rangle_t}{\langle I \rangle_i}$$

10.2.15 Composite surface transmission coefficient

$$\alpha_{t_{\text{eff}}} = \frac{\sum_{j=1}^n \alpha_{t_j} A_j}{A_T} \quad \text{where } A \text{ is surface area}$$

10.2.16 Transmission loss

$$\begin{aligned} \text{T.L.} &= 10 \log \left(\frac{1}{\alpha_t} \right) \\ &= 10 \log \left[\frac{(\sigma\omega)^2}{(2\rho_0 a_0)^2} \right] + 10 \log \left[1 + \frac{(2\rho_0 a_0)^2}{(\sigma\omega)^2} \right] \end{aligned}$$

10.2.17 Spherical wave equation

$$p = \frac{A}{r} e^{i(\omega t - kr)}$$

10.2.18 Spherical Sound Pressure Level

$$\begin{aligned} \text{SPL} &= 20 \log \left(\frac{p_{\text{rms}}}{p_{\text{ref}}} \right) \\ &= 20 \log \left(\frac{A}{\sqrt{2} p_{\text{ref}}} \right) - 20 \log(r) \\ \Delta_{\text{SPL}} &= \text{SPL}_1 - \text{SPL}_2 = 20 \log \left(\frac{r_2}{r_1} \right) \end{aligned}$$

10.2.19 Spherical acoustic velocity

$$\begin{aligned} u'_r &= \frac{p'(t)}{\rho_0 a_0} \left(1 + \frac{1}{ikr} \right) \\ &= \frac{A}{\rho_0 a_0 r} \left(1 + \frac{1}{ikr} \right) e^{i(\omega t - kr)} \end{aligned}$$

10.2.20 Reverberant Field energy density

$$\begin{aligned} \langle e \rangle_{\text{RF}} &= \frac{4\pi p_{\text{rms}}^2}{\rho_0 a_0^2} \\ &= \frac{2\pi A^2}{\rho_0 a_0^2} \end{aligned}$$

10.2.21 Reverberant field intensity

$$\begin{aligned} \langle I \rangle_{\text{RF}} &= \frac{\pi A^2}{2\rho_0 a_0} \\ &= \frac{\langle e \rangle_{\text{RF}} a_0}{4} \end{aligned}$$

10.2.22 Absorption

$$\begin{aligned} \langle I \rangle_{\text{RF}} &= \frac{\mathbb{P}(1 - \alpha_A)}{\alpha_A A_{\text{wall}}} = \frac{\mathbb{P}}{R} \\ \alpha_A &= \frac{\langle I \rangle_a}{\langle I \rangle_i} = \frac{\langle I \rangle_i - \langle I \rangle_r}{\langle I \rangle_i} \end{aligned}$$

10.2.23 Reverberant Field SPL

$$\begin{aligned} \text{SPL}_{\text{RF}} &= 10 \log_{10} \left[\frac{(p_{\text{rms}}^2)_{\text{RF}}}{p_{\text{ref}}^2} \right] \\ &= 10 \log \left(\frac{4\rho_0 a_0 \mathbb{P}}{R p_{\text{ref}}^2} \right) \\ R &= \frac{A_{\text{wall}} \alpha_A}{(1 - \alpha_A)} \end{aligned}$$

10.2.24 Direct field and Total SPL

$$\text{SPL}_{\text{total}} = 10 \log_{10} \left(\frac{\rho_0 a_0 \mathbb{P}}{p_{\text{ref}}^2} \right) + 10 \log_{10} \left(\frac{1}{4\pi r^2} + \frac{4}{R} \right)$$

10.2.25 Average absorption coefficient

$$\alpha_A = \frac{\sum \alpha_{A_i} A_i}{\sum A_i}$$

10.2.26 Average SPL or "equivalent level"

$$\text{SPL}_{\text{AVG}} = L_{\text{EQ}} = 10 \log_{10} \left(\frac{1}{T} \sum_{i=1}^N 10^{\text{SPL}_i/10} \Delta t_i \right)$$

10.2.27 Day/Night level

$$L_{\text{DN}} = 10 \log_{10} \left[\frac{1}{T} \left(\sum_{0700}^{2200} 10^{\text{SPL}(t)/10} \Delta t_i + \sum_{2200}^{0700} 10^{\frac{\text{SPL}(t)+10}{10}} \Delta t_i \right) \right]$$

10.2.28 Permissible time & dose

$$T = \frac{8}{2^{\frac{L-90}{5}}}$$
$$D = \sum_{i=1}^N \frac{A_i}{T} = \frac{1}{8} \sum_{i=1}^N \Delta t_i 2^{\frac{L_i-90}{5}} \leq 1$$

A_i is the actual time

T is the permissible time (8 hours at 90 dB)

11 Numerical Analysis

11.1 Second Order PDEs classification

11.1.1 General Form

$$A \frac{\partial^2 \phi}{\partial x^2} + B \frac{\partial^2 \phi}{\partial x \partial y} + C \frac{\partial^2 \phi}{\partial y^2} + D \frac{\partial \phi}{\partial x} + E \frac{\partial \phi}{\partial y} + f\phi + G = 0$$

$$A \frac{\partial^2 \phi}{\partial x^2} + B \frac{\partial^2 \phi}{\partial x \partial y} + C \frac{\partial^2 \phi}{\partial y^2} = H$$

$$H = - \left(D \frac{\partial \phi}{\partial x} + E \frac{\partial \phi}{\partial y} + f\phi + G \right)$$

11.1.2 Characteristics in physical space

$$\left(\frac{dy}{dx} \right)_{\alpha, \beta} = \frac{B \pm \sqrt{B^2 - 4AC}}{2A}$$

11.1.3 Equation characteristics

elliptic if	$B^2 - 4AC < 0$
parabolic if	$B^2 - 4AC = 0$
hyperbolic if	$B^2 - 4AC > 0$

11.2 Finite Difference Schemes

11.2.1 First derivative

$$u'_i = \frac{1}{\Delta x} (-u_{i-1} + u_i) + \frac{1}{2} \Delta x u''$$

$$= \frac{1}{\Delta x} (-u_i + u_{i+1}) - \frac{1}{2} \Delta x u''_i$$

$$= \frac{1}{\Delta x} (-u_{i-1} + u_{i+1}) - \frac{1}{6} \Delta x^2 u_i^{(3)}$$

$$= \frac{1}{\Delta x} \left(\frac{1}{2} u_{i-2} - 2u_{i-1} + \frac{3}{2} u_i \right) + \frac{1}{3} \Delta x^2 u_i^{(3)}$$

$$= \frac{1}{\Delta x} \left(-\frac{3}{2} u_i + 2u_{i+1} - \frac{1}{2} u_{i+2} \right) + \frac{1}{3} \Delta x^2 u_i^{(3)}$$

11.3 Stability Analysis

11.3.1 Discrete Perturbation Stability analysis

Consider the parabolic model equation

$$\frac{\partial u}{\partial t} = \alpha \frac{\partial^2 u}{\partial x^2}$$

Using a first order time and second order spatial derivative, this equation may be written as:

$$\frac{u_i^{n+1} - u_i^n}{\Delta t} = \alpha \frac{u_{i+1}^n - 2u_i^n + u_{i-1}^n}{(\Delta x)^2}$$

If a disturbance ϵ at node i and time level n is introduced and we search for solution at time level $n + 1$ for all i nodes, the finite difference eq. becomes:

$$\frac{u_i^{n+1} - (u_i^n + \epsilon)}{\Delta t} = \alpha \frac{u_{i+1}^n - 2(u_i^n + \epsilon) + u_{i-1}^n}{(\Delta x)^2}$$

If $u^n = 0$ for all i the equation reduces to:

$$\frac{u_i^{n+1} - \epsilon}{\Delta t} = \alpha \frac{-2\epsilon}{(\Delta x)^2}$$

or equivalently

$$u_i^{n+1} = \epsilon \left\{ 1 + 2\alpha \left[\frac{\Delta t}{(\Delta x)^2} \right] \right\}$$

$$= \epsilon (1 - 2d)$$

$$\rightarrow \frac{u_i^{n+1}}{\epsilon} = (1 - 2d)$$

where $d = \alpha \Delta t / (\Delta x)^2$ is known as the *diffusion number*. It's required that,

$$\left| \frac{u_i^{n+1}}{\epsilon} \right| \leq 1$$

or

$$1 - 2d \leq 1 \quad \text{and} \quad 1 - 2d \geq -1$$

11.3.2 Von Neumann Stability Analysis

Solution of finite difference equation is expanded in a Fourier series. The decay or growth of the *amplification factor* determines the stability.

Assume a Fourier component for u_i^n as

$$u_i^n = U^n e^{IP(\Delta x)i}$$

where $I = \sqrt{-1}$, U^n is the amplitude at time n , and P is the wave number in the x -direction, i.e., $\lambda_x = 2\pi/P$, where λ_x is the wavelength. Similarly,

$$u_i^{n+1} = U^{n+1} e^{IP(\Delta x)i} \quad \text{and} \quad u_{i\pm 1}^n = U^n e^{IP(\Delta x)(i\pm 1)}$$

If phase angle $\theta = P\Delta x$ is defined, then

$$u_i^n = U^n e^{I\theta i}$$

$$u_i^{n+1} = U^{n+1} e^{I\theta i}$$

$$u_{i\pm 1}^n = U^n e^{I\theta(i\pm 1)}$$

Consider again

$$\frac{u_i^{n+1} - u_i^n}{\Delta t} = \alpha \frac{u_{i+1}^n - 2u_i^n + u_{i-1}^n}{(\Delta x)^2}$$

or in terms of the diffusion number

$$u_i^{n+1} = u_i^n + d(u_{i+1}^n - 2u_i^n + u_{i-1}^n)$$

Substituting the Fourier component and cancelling out terms of $e^{I\theta i}$, we get

$$U^{n+1} = U^n [1 + d(e^{I\theta} - e^{-I\theta} - 2)]$$

Using the relation $\cos \theta = (e^{I\theta} + e^{-I\theta})/2$, we get

$$U^{n+1} = U^n [1 + 2d(\cos \theta - 1)]$$

Introducing the amplification factor $U^{n+1} = GU^n$, we get

$$G = 1 - 2d(1 - \cos \theta)$$

Stability requires,

$$|G| \leq 1 \quad \text{and} \quad |1 - 2d(1 - \cos \theta)| \leq 1$$

so that

$$\begin{aligned} 1 - 2d(1 - \cos \theta) &\leq 1 \\ &\text{and} \\ 1 - 2d(1 - \cos \theta) &\geq -1 \end{aligned}$$

must be valid for all values of θ . With a maximum value of $(1 - \cos \theta) = 2$, the LHS becomes $1 - 4d$. Solving the new equation gives us a final condition

$$d \leq \frac{1}{2} \quad (3)$$

11.3.3 Scheme Requirements

- *Consistency*: A finite difference approximation of a PDE is consistent if the finite difference equation approaches the PDE as the grid size approaches zero.
- *Stability*: A numerical scheme is said to be stable if any error introduced in the finite difference equation does not grow with the solution of the finite difference equation.
- *Convergence*: A finite difference scheme is convergent if the solution of the finite difference scheme approaches that of the PDE as the grid size approaches zero.

- *Lax's equivalence theorem*: For a FDE which approximates a well-posed, linear initial value problem, the necessary and sufficient condition for convergence is that the FDE must be stable and consistent.
- *The conservative (divergent) form of a PDE*: In this formulation of a physical law, the coefficients of the derivatives are either constant or, if variable, their derivatives do not appear anywhere in the equation.
- *Conservative property of a FDE*: If the finite difference approximation of a PDE maintains the integral property of the conservation law over an arbitrary region containing any number of grid points, it is said to possess a conservative property.

11.4 Classification of equations

Elliptic:	$\frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} = 0$
Parabolic:	$\frac{\partial T}{\partial t} = \alpha \frac{\partial^2 T}{\partial x^2}$
Hyperbolic:	$\frac{\partial^2 \phi}{\partial t^2} = a^2 \frac{\partial^2 \phi}{\partial x^2}$

12 Useful Identities

12.1 Series

12.1.1 Expansion of asymptotic series to a power

$$\begin{aligned} (x_0 + \epsilon x_1 + \epsilon^2 x_2 + \epsilon^3 x_3 + \dots)^n &= x_0^n \\ &+ \epsilon n x_0^{n-1} x_1 \\ &+ \epsilon^2 \left[\frac{n(n-1)}{2!} x_0^{n-2} x_1^2 + n x_0^{n-1} x_2 \right] \\ &+ \epsilon^3 \left[\frac{n(n-1)(n-2)}{3!} x_0^{n-3} x_1^3 + n(n-1) x_0^{n-2} x_1 x_2 + n x_0^{n-1} x_3 \right] \\ &+ \epsilon^4 \left[\frac{n(n-1)(n-2)(n-3)}{4!} x_0^{n-4} x_1^4 + \frac{n(n-1)(n-2)}{2!} x_0^{n-3} x_1^2 x_2 + \frac{n(n-1)}{2!} x_0^{n-2} (x_2^2 + 2x_1 x_3) + n x_0^{n-1} x_4 \right] \end{aligned}$$